

# On using convolutions with exponential distributions for solving a Kolmogorov backward equation

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# Introduction

- 1 Problem setup
- 2 Randomization and approximation
- 3 Finding the solution in terms of EPV operators  $\mathcal{E}^{\pm}$
- 4 Calculating convolutions
- 5 References

# Contents

- 1 Problem setup
- 2 Randomization and approximation
- 3 Finding the solution in terms of EPV operators  $\mathcal{E}^\pm$
- 4 Calculating convolutions
- 5 References

## A diffusion equation

A 3-dimensional partial differential equation for function  $u = u(x, t)$ , where  $x = (x_1, x_2) \in \mathbb{R}^2$ ,  $t \in \mathbb{R}$ :

General form (Øksendal)

$$\left(\frac{\partial}{\partial t} + L\right)u = 0, \quad L = \sum_i \mu_i(x) \frac{\partial}{\partial x_i} + \frac{1}{2} \sum_{i,j} (\sigma \sigma^T)_{ij}(x) \frac{\partial^2}{\partial x_i \partial x_j},$$

$i = 1, 2, j = 1, 2$ ; functions  $\mu(x) = (\mu_1(x), \mu_2(x))$  and

$$\sigma(x) = \begin{pmatrix} \sigma_{11}(x) & \sigma_{12}(x) \\ \sigma_{21}(x) & \sigma_{22}(x) \end{pmatrix}$$

satisfies conditions of Theorem 5.2.1.  $\mu : \mathbf{R}^2 \rightarrow \mathbf{R}^2$  — drift coefficient  
 $\sigma : \mathbf{R}^2 \rightarrow \mathbf{R}^2 \times \mathbf{R}^2$  (or  $\frac{1}{2}(\sigma \sigma^T)$ ) — diffusion coefficient

- With a suitable initial and boundary conditions, the Kolmogorov backward equation can be solved (Øksendal; Eström (2010)).
- There are no general analytical solutions.
- Numerical methods are based on the exact form of  $\mu_i(x)$  and  $\sigma(x)$ .

Let  $T > 0$  be a time moment,  $H > 0$  – an absorbing barrier,  $g(x) : \mathbf{R} \rightarrow \mathbf{R}^{\geq 0}$  – some suitable function, which decays rapidly on infinity.

### Terminal and boundary conditions

$$\begin{cases} \left( \frac{\partial}{\partial t} + L \right) u = 0, & x_1 > H, t < T, \\ u(x_1, x_2, T) = g(x_1), & x_1 > H, \\ u(x_1, x_2, t) = 0, & x_1 \leq H, t \leq T. \end{cases}$$

# A diffusion equation

- The solution to a diffusion equation in a given domain can be interpreted as an expectation (Dynkin, Feinman, Katz).
- The idea was later generalized for the case of Lévy processes.

# Ito diffusions - general

## Two correlated diffusions

$$\begin{cases} dX_1(t) = \mu_1 dt + \sigma_{11} dB_1(t) + \sigma_{12} dB_2(t), \\ dX_2(t) = \mu_2 dt + \sigma_{21} dB_1(t) + \sigma_{22} dB_2(t). \end{cases}$$

- $B_1(t), B_2(t)$  are Brownian motions (Wiener processes)

## An example from mathematical finance

Define const.  $\kappa_V > 0, \theta_V > 0, \sigma_V > 0$ ;  $\rho \in (-1, 1), \hat{\rho} = \sqrt{1 - \rho^2}$ . Denote  $X_1(t)$  as  $S_t$  and  $X_2(t)$  as  $V_t$ . Assume  $\mu_1 = 0$ ,  $\mu_2 = \kappa_V(\theta_V - V_t)$  and

$$\sigma = \begin{pmatrix} \rho\sqrt{V_t}S_t & \hat{\rho}\sqrt{V_t}S_t \\ \sigma_V\sqrt{V_t} & 0 \end{pmatrix},$$

Processes:

$$\begin{cases} dS_t = \sqrt{V_t}S_t(\rho dB_1(t) + \hat{\rho}dB_2(t)), \\ dV_t = \kappa_V(\theta_V - V_t)dt + \sigma_V\sqrt{V_t}dB_1(t). \end{cases} \quad (1)$$

An infinitesimal operator  $L$ :

$$L = \frac{1}{2}S^2v\frac{\partial^2}{\partial S^2} + \rho\sigma_VvS\frac{\partial^2}{\partial S\partial v} + \frac{1}{2}\sigma_V^2v\frac{\partial^2}{\partial v^2} + \kappa_V(\theta_V - v)\frac{\partial}{\partial v}. \quad (2)$$



## The problem in $u(S, v, t)$ terms

$$\left\{ \begin{array}{l} \left( \frac{\partial}{\partial t} + \frac{1}{2} S^2 v \frac{\partial^2}{\partial S^2} + \rho \sigma_V v S \frac{\partial^2}{\partial S \partial v} + \right. \\ \left. + \frac{1}{2} \sigma_V^2 v \frac{\partial^2}{\partial v^2} + \kappa_V (\theta_V - v) \frac{\partial}{\partial v} \right) u = 0, \quad S > H, v > 0, t < T, \\ \\ u(S, v, T) = g(S), \quad S > H, v > 0, \\ u(S, v, t) = 0, \quad S \leq H, v > 0, t \leq T. \end{array} \right.$$

The solution (exist. & uniq. — Cont, Tankov (2004)):

$$u(S, v, 0) = E[\mathbf{1}_{(T, +\infty)}(T_H) \cdot g(S_T) | S_0 = S, V_0 = v], \quad t = 0. \quad (3)$$

## The substitution (Zanette, Briani et al. (2017))

The structure:

$$Y_t = \ln\left(\frac{S_t}{H}\right) - \frac{\rho}{\sigma_V} V_t, \quad S_t = H \exp\left(Y_t + \frac{\rho}{\sigma_V} V_t\right).$$

The system in its terms:

$$\begin{cases} dY_t = \mu_Y(V_t)dt + \hat{\rho}\sqrt{V_t}dB_2(t), \\ dV_t = \mu_V(V_t)dt + \sigma_V\sqrt{V_t}dB_1(t), \end{cases}$$

where

$$\mu_Y(v) = -\frac{1}{2}v - \frac{\rho}{\sigma_V}\kappa_V(\theta_V - v)$$

and

$$\mu_V(v) = \kappa_V(\theta_V - v).$$

With  $g(y) = g(He^{y + \frac{\rho}{\sigma_V} v})$ :

The problem in terms of  $f(y, v, t) := u\left(H \exp(y + \frac{\rho}{\sigma_V} v), v, t\right)$

$$\begin{cases} \left( \frac{\partial}{\partial t} + \frac{1}{2} \hat{\rho}^2 v \frac{\partial^2}{\partial y^2} + \frac{1}{2} \sigma_V^2 v \frac{\partial^2}{\partial v^2} + \mu_Y(v) \frac{\partial}{\partial y} + \mu_V(v) \frac{\partial}{\partial v} \right) f = 0, \\ y + \frac{\rho}{\sigma_V} v > 0, v > 0, t < T, \\ f(y, v, T) = g(y), & y + \frac{\rho}{\sigma_V} v > 0, v > 0, \\ f(y, v, t) = 0, & y + \frac{\rho}{\sigma_V} v \leq 0, v > 0, t \leq T. \end{cases}$$

$T_H$  is as the earliest time the process  $Y_t + \frac{\rho}{\sigma} V_t$  enters  $(-\infty, 0]$ :

$$T_H = \inf_{t \geq 0} \left\{ t : Y_t + \frac{\rho}{\sigma_V} V_t \leq 0 \right\}.$$

The expectation

$$f(y, v, 0) = E^{y, v} [\mathbf{1}_{(T, +\infty)}(T_H) g(Y_T)].$$

# Contents

- 1 Problem setup
- 2 Randomization and approximation**
- 3 Finding the solution in terms of EPV operators  $\mathcal{E}^\pm$
- 4 Calculating convolutions
- 5 References

## Carr randomization

Let  $N \in \mathbb{N}$ ,  $\Delta t = T/N$ ,  $q > 0$ ,  $T_q \sim \text{Exp}(\Delta t^{-1})$ .

We calculate a sequence of

$$f_n(y, v) \approx f(y, v, \sum_{j=1}^n T_q^j), \quad n = 0, 1, \dots, N;$$

where  $f_0(y, v) \approx f(y, v, 0)$ ;  $f_N(y, v) = g(y)$ ,  
and  $T_q^j \sim \text{Exp}(\Delta t^{-1})$  is a sequence of ind. random variables.

### Expectations

$$f_n(y) = E^{y,v} \left[ \mathbf{1}_{T_q^{n+1}, +\infty} \left( \inf_{t \geq 0} \left\{ t : Y_{T_q^{n+1}} + \frac{\rho}{\sigma_V} V_{T_q^{n+1}} \leq 0 \right\} \right) \cdot f_{n+1}(Y_{T_q^{n+1}}, V_{T_q^{n+1}}) \right]$$

# The Markov chain

## The recombining binomial tree:

$$V_{n,k} = \left( \sqrt{V_0} + \frac{\sigma_V}{2} (2k - n) \sqrt{\Delta t} \right)^2 \cdot \mathbf{1}_{(0,+\infty)} \left( \sqrt{V_0} + \frac{\sigma_V}{2} (2k - n) \sqrt{\Delta t} \right),$$

where  $n = 0, 1, \dots, N$ ,  $k = 0, 1, \dots, n$ .

## Transitions

From  $(n, k)$  to either  $(n + 1, k_u)$  or  $(n + 1, k_d)$ .

With probabilities  $p_u$  and  $p_d$ .

The values  $k_u, k_d$  and  $p_u, p_d$  are based  $V_t$ , like in (Briani, Zanette et al (2013)).

# Expectations

Recurrent calculation w.r.t the Markov chain:

$$f_{n,k}(y) = \mathbf{1}_{(0,+\infty)}\left(y + \frac{\rho}{\sigma_V} V_{n,k}\right) \cdot (p_u f_{n,k}^u(y) + p_d f_{n,k}^d(y)),$$

where

$$f_{n,k}^u(y) = E^y \left[ \mathbf{1}_{(0,+\infty)}\left(Y_{T_q}^{n+1,k_u} + \frac{\rho}{\sigma_V} V_{n+1,k_u}\right) \cdot f_{n+1}\left(Y_{T_q}^{n+1,k_u}, V_{n+1,k_u}\right) \right],$$

and  $Y_{T_q}^{n,k}$  is a value of  $Y_{T_q}$  from the system, when  $V_{T_q} = V_{n,k}$ ;

the value of  $f_{n,k}^d(y)$  is calculated analogously.

# Contents

- 1 Problem setup
- 2 Randomization and approximation
- 3 Finding the solution in terms of EPV operators  $\mathcal{E}^\pm$**
- 4 Calculating convolutions
- 5 References



## Expected Present Value (S. Levendorskii) operator

$$\mathcal{E}_q F(x) = qE^x \left[ \int_0^{+\infty} e^{-qt} F(X_t) dt \right] = E[F(x + X_{T_q})],$$

where  $X_t$  is an Ito diffusion.

## Sup and inf processes with the EPVs

$$\overline{X}_t = \sup_{0 \leq s \leq t} X_s, \quad \underline{X}_t = \inf_{0 \leq s \leq t} X_s,$$

$$\mathcal{E}_q^+ F(x) = qE^x \left[ \int_0^{+\infty} e^{-qt} F(\overline{X}_t) dt \right] = E[F(x + \overline{X}_{T_q})],$$

$$\mathcal{E}_q^- F(x) = qE^x \left[ \int_0^{+\infty} e^{-qt} F(\underline{X}_t) dt \right] = E[F(x + \underline{X}_{T_q})].$$

## EPV in the problem considered

Using the property  $\bar{X}_t \stackrel{d}{\sim} X_t - \underline{X}_t$ , one can quite easily obtain:

The solution in terms of the EPV operators

$$f_{n,k}^u(y) = q\mathcal{E}_q^- \left( \mathbf{1}_{(0,+\infty)}(y + \frac{\rho}{\sigma_V} V_{n+1,k_u}) \cdot \mathcal{E}_q^+ f_{n+1,k_u}(y) \right),$$

and use an analogous formula for  $f_{n,k}^d(y)$ .

EPV operators as convolutions

$$\mathcal{E}_q F(x) = \int_{-\infty}^{+\infty} F(x+u) P_q(du), \quad \mathcal{E}_q^\pm F(x) = \int_{-\infty}^{+\infty} F(x+u) P_q^\pm(du),$$

where  $P_q^+(-\infty, 0) = 0$ ,  $P_q^-(0, +\infty) = 0$ .

## Wiener-Hopf factorization for a symbol of a PDO

If  $X_t$  is a Lévy process, and its char. exp. is a function  $\psi(\xi)$ , defined by the Levy-Khintchine formula, then the char. functions for the distributions  $P_q^\pm(du)$  are  $\phi^\pm(\xi)$ , such that:

$$\phi^+(\xi) \cdot \phi^-(\xi) = q(q + \psi(\xi))^{-1}. \quad (4)$$

There also exist constants  $\omega_- < 0 < \omega_+$  such that  $\phi^+(\xi)$  admits an analytical cont. into half-plane  $\text{Im } \xi > \omega_-$  (and  $\phi^-(\xi)$  — into  $\text{Im } \xi < \omega_+$ ).

### PDO representation and symbols

$q(q + \psi(\xi))^{-1}$  is a symbol of  $\mathcal{E}_q$   
 $\phi^\pm(\xi)$  are symbols of  $\mathcal{E}_q^\pm$ .

# Contents

- 1 Problem setup
- 2 Randomization and approximation
- 3 Finding the solution in terms of EPV operators  $\mathcal{E}^\pm$
- 4 Calculating convolutions**
- 5 References

The characteristic function of  $Y_t^{n,k}$ ,  $t < T_q$ , is

$$\psi(\xi) = \frac{\sigma_{n,k}^2}{2} \xi^2 - i\gamma_{n,k}\xi,$$

where  $\sigma_{n,k} = \hat{\rho}\sqrt{V_{n,k}}$ ,  $\gamma_{n,k} = \mu_Y(V_{n,k})$ . For this case:

$$\phi^+(\xi) = \frac{\beta_q^+}{\beta_q^+ - i\xi}, \quad \phi^-(\xi) = \frac{-\beta_q^-}{-\beta_q^- + i\xi},$$

where

$$\beta_q^+ = \frac{-\gamma_{n,k} + \sqrt{\gamma_{n,k}^2 + 2\sigma_{n,k}^2 q}}{\sigma_{n,k}^2}, \quad \beta_q^- = \frac{-\gamma_{n,k} - \sqrt{\gamma_{n,k}^2 + 2\sigma_{n,k}^2 q}}{\sigma_{n,k}^2}.$$

The respective distributions:

$$P_q^-(du) = -\beta_q^- e^{-\beta_q^- u} \mathbf{1}_{(-\infty, 0)}(u) du, \quad P_q^+(du) = \beta_q^+ e^{-\beta_q^+ u} \mathbf{1}_{(0, +\infty)}(u) du$$

PDO representation ( $O n \log(n)$  operations via FFT):

$$\varepsilon_q^\pm u(x) = (2\pi)^{-1} \int_{-\infty}^{+\infty} e^{ix\xi} \phi_q^\pm(\xi) \hat{u}(\xi) d\xi.$$

Let us define a reasonably dense grid  $y_k = y^* + yh$ ,  $k = 0, 1, \dots, M$ , where  $M \in \mathbb{N}$  is large, denote  $F_n(y) := f_{n,k}^u(y)$ ,  $F_{n+1}(y) := f_{n+1,k_u}^u(y)$  and consider the equation:

Convolutions calculation

$$\mathcal{E}^+ F_{n+1}(y_k) = \int_0^{+\infty} \beta_q^+ e^{-\beta_q^+ u} F_{n+1}(y_k + u) du. \quad (5)$$

## $\mathcal{E}^+$ , step 1

As  $g(y)$  decays on infinity,  $\lim_{y \rightarrow +\infty} F_n(y) = \lim_{y \rightarrow +\infty} g(y) = 0$ , we restrict the integration area with  $y^* > 0$ . If  $y_k + u < y^*$ , then  $u < y^* - y_k$ . We get:

$$\mathcal{E}^+ F_{n+1}(y_k) \approx \int_0^{y^* - y_k} \beta_q^+ e^{-\beta_q^+ u} F_{n+1}(y_k + u) du.$$

Make a substitution:  $w = u + y_k$ :

$$\mathcal{E}^+ F_{n+1}(y_k) \approx e^{\beta_q^+ y_k} \int_{y_k}^{y^*} \beta_q^+ e^{-\beta_q^+ w} F_{n+1}(w) dw.$$

## $\mathcal{E}^+$ , step 2

From an indicator: for  $y_k - h < -\frac{\rho}{\sigma_V} V_{n+1, k_u}$ ,  $F_{n+1}(y_k - h) = 0$ .

$$\mathcal{E}^+ F_{n+1}(y_{k-1}) \approx e^{-\beta_q^+ h} \mathcal{E}^+ F_{n+1}(y_k) + e^{-\beta_q^+ h} \int_{y_k - h}^{y_k} \beta_q^+ e^{-\beta_q^+ (w - y_k)} F_{n+1}(w) dw.$$

Trapezoid approximation for the integral part:

$$\frac{h}{2} \cdot \beta_q^+ \left( F_{n+1}(y_{k-1}) + e^{-\beta_q^+ h} F_{n+1}(y_k) \right)$$

Simpson approximation example ( $\mathcal{E}^+ F_{n+1}(y_{k-2})$  used):

$$\frac{h}{3} \cdot \beta_q^+ \left( F_{n+1}(y_{k-2}) + 4e^{-\beta_q^+ h} F_{n+1}(y_{k-1}) + e^{-2\beta_q^+ h} F_{n+1}(y_k) \right)$$



## $\mathcal{E}^-$ , step 1

Let us put  $z = \frac{\rho}{\sigma_V} V_{n+1, k_u}$  and consider:

$$\begin{aligned} q\mathcal{E}^- \left( \mathbf{1}_{(0, +\infty)}(y_k + z) \cdot \mathcal{E}^+ F_{n+1}(y_k) \right) &= \\ &= q \int_{-(z+y_k)}^0 -\beta_q^- e^{-\beta_q^- u} \cdot \mathcal{E}^+ F_{n+1}(u + y_k) du \end{aligned}$$

The integral form for  $q\mathcal{E}^-$  in  $y_k$ :

$$\stackrel{w:=u+y_k}{=} q e^{\beta_q^- y_k} \int_{-z}^{y_k} -\beta_q^- e^{-\beta_q^- w} \mathcal{E}^+ F_{n+1}(w) dw.$$

## $\mathcal{E}^-$ , step 2

Then, for  $y_{k+1} = y_k + h$ , using the same reasoning we get:

$$\begin{aligned} q\mathcal{E}^-(\mathbf{1}_{(0,+\infty)}(y_{k+1} + z) \cdot \mathcal{E}^+ F_{n+1}(y_{k+1})) &= \\ &= qe^{\beta_q^-(y_k+h)} \left( e^{\beta_q^- h} \mathcal{E}^- \mathbf{1}_{(-\infty,0)}(y_k + z) + \right. \\ &\quad \left. + \int_{y_k}^{y_k+h} -\beta_q^- e^{-\beta_q^- w} \mathcal{E}^+ F_{n+1}(w) dw \right) \end{aligned}$$

### Trapezoid approximation

$$\frac{h}{2} \cdot (-\beta_q^-) \left( \mathcal{E}^+ F_{n+1}(y_{k+1}) + e^{\beta_q^- h} \mathcal{E}^+ F_{n+1}(y_k) \right)$$

### Simpson approximation ( $\mathcal{E}^- F_{n+1}(y_{k+2})$ used)

$$\frac{h}{3} \cdot (-\beta_q^-) \left( e^{2\beta_q^- h} \mathcal{E}^+ F_{n+1}(y_{k+2}) + e^{\beta_q^- h} \mathcal{E}^+ F_{n+1}(y_{k+1}) + \mathcal{E}^+ F_{n+1}(y_k) \right)$$

# Summary

- Using the procedure above to solve the arising 1-dimensional problems, we obtain an approximate solution for the Kolmogorov backward equation.
- The scheme presented is an iterative method for calculating integrals, where already calculated values of  $\mathcal{E}^\pm F_{n+1}(y_k)$  are used to obtain values for the neighboring  $y_k$ .

## The main advantage

We only need  $O(n)$  operations to calculate the integrals

THANK YOU FOR ATTENTION!

# Contents

- 1 Problem setup
- 2 Randomization and approximation
- 3 Finding the solution in terms of EPV operators  $\mathcal{E}^\pm$
- 4 Calculating convolutions
- 5 References

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